We know (from Fourier-series theory) that

$$\rho_k = \frac{1}{2\pi} \int_0^{2\pi} \omega(\beta) \cos(k\beta) d\beta$$

Introducing $y \equiv \exp(i\beta)$, which implies that

$$\cos(k\beta) = \frac{y^k + y^{-k}}{2}$$
 and $d\beta = \frac{dy}{iy}$

this can be converted to

$$\frac{1}{2\pi i} \oint \omega(y) \left(\frac{y^{k-1} + y^{-k-1}}{2}\right) dy \tag{1}$$

where the integration is over the *unit circle* centered at 0 (counterclockwise). The theory of *contour integration* tells us that the result is the sum of *residues* of the integrand at all *singularities* inside this circle. Note that the ω (being, originally, a function of

$$\cos\beta = \frac{y + \frac{1}{y}}{2}$$

has the $\omega(y) = \omega(\frac{1}{y})$ symmetry.

Expanding $\omega(y)$, which is a rational function of y, in terms of partial fractions makes it quite easy to identify the singularities, and to find their residues. One should realize that each term of this expansion will have the form of either

$$c \cdot y^m \tag{2}$$

where c is a constant (different for different terms) and m is an integer (positive, negative, or zero), or of

$$\frac{c}{(y-\lambda)^{\ell}}\tag{3}$$

where c and $\lambda \neq 0$ are two constants (potentially complex), and ℓ is a positive integer (usually equal to 1). Clearly any such term with complex c and λ will have its *complex conjugate* counterpart. But, more importantly, due to the $y \leftrightarrow \frac{1}{y}$ symmetry, each term of the type $c \cdot y^m$ must be 'paired' with the corresponding $c \cdot y^{-m}$ term (this time, with the *same* c), and similarly every term of type (3) - regardless whether c and λ are complex or real - will have its 'dual' set of terms, given by the partial-fraction expansion of

$$\frac{c}{\left(\frac{1}{y}-\lambda\right)^{\ell}} = \frac{cy^{\ell}}{(1-y\lambda)^{\ell}} = \frac{c\left(-\frac{1}{\lambda}\right)^{\ell}y^{\ell}}{\left(y-\frac{1}{\lambda}\right)^{\ell}}$$
(4)

namely

$$c\left(-\frac{1}{\lambda}\right)^{\ell}\sum_{j=0}^{\ell}\binom{\ell}{j}\frac{\lambda^{-j}}{\left(y-\frac{1}{\lambda}\right)^{j}}$$

(the details of which become irrelevant shortly). Note that the absolute value of each λ must be different from 1 (to represent a stationary model).

One can prove that the contributions of cy^m and cy^{-m} to (1) are identical (for all positive-integer values of k); the same is true for the (3) and (4) pair (see below). This means that we can divide terms of the partial-fraction expansion of $\omega(y)$ into two parts:

- 1. those which have the form of either (2) with *m* negative or of (3) with $|\lambda| < 1$,
- 2. all their dual terms these consist of positive powers of y, and terms or Type (3) with $|\lambda| > 1$, all easily identifiable.

The constant term is an exception to this, but since it contributes only to $\rho_0 \equiv 1$, we can ignore it. We have excluded the $|\lambda| = 1$ possibility which cannot happen, as explained earlier.

All we have to do in the end is to find the sum of residues of terms of Type 1 and multiply the result by 2 (totally ignoring terms of Type 2, which would contribute the same amount - that means we never need to utilize the last two formulas). Note that multiplying by 2 is the same as replacing

$$\left(\frac{y^{k-1}+y^{-k-1}}{2}\right)$$

in (1) by $(y^{k-1} + y^{-k-1})$ - that's what we will do from now on. **Proof.** First we prove that

$$\frac{1}{2\pi i} \oint y^m \left(y^{k-1} + y^{-k-1} \right) dy = \frac{1}{2\pi i} \oint y^{-m} \left(y^{k-1} + y^{-k-1} \right) dy = \delta_{m,k}$$

when m is a positive integer. Note that $\delta_{m,k}$ is the Kronecker's delta (equal to 1 when m = k, equal to 0 otherwise).

The result is quite obvious by realizing that the only power of y which integrates (in the above sense) to $2\pi i$ is y^{-1} - all the other integer powers yield zero. It is thus the y^{m-k-1} part of the first integral and the y^{-m+k-1} part of the second integral which result in $\delta_{m,k}$ (the y^{m+k-1} and y^{-m-k-1} terms contribute 0).

Secondly, we have to deal with the

$$\frac{1}{2\pi i} \oint (y-\lambda)^{-\ell} \left(y^{k-1} + y^{-k-1} \right) dy = \frac{1}{2\pi i} \oint \frac{\left(-\frac{1}{\lambda}\right)^{\ell} y^{\ell}}{\left(y-\frac{1}{\lambda}\right)^{\ell}} \left(y^{k-1} + y^{-k-1} \right) dy$$

identity, assuming that $|\lambda| = 1$. The LHS integral has two singularities inside the unit circle: one at 0 and the other at λ . The first one contributes (another well-known result of contour integration):

$$\frac{1}{k!} \cdot \frac{d^k (y-\lambda)^{-\ell}}{dy^k} \bigg|_{y=0} = \frac{(-\ell)(-\ell-1)\dots(-\ell-k+1)}{k!} \cdot (-\lambda)^{-\ell-k} = (-1)^\ell \binom{\ell+k-1}{k} \lambda^{-\ell-k}$$

the second one adds

$$\begin{aligned} \frac{1}{(\ell-1)!} \cdot \frac{d^k y^{k-1}}{dy^{\ell-1}} \bigg|_{y=\lambda} &+ \frac{1}{(\ell-1)!} \cdot \frac{d^k y^{-k-1}}{dy^{\ell-1}} \bigg|_{y=\lambda} = \\ \frac{(k-1)(k-2)\dots(k-\ell+1)}{(\ell-1)!} \cdot \lambda^{k-\ell} &+ \frac{(-k-1)(-k-2)\dots(-k-\ell+1)}{(\ell-1)!} \cdot \lambda^{-k-\ell} = \\ \binom{k-1}{\ell-1} \lambda^{k-\ell} &+ (-1)^{\ell-1} \binom{\ell+k-1}{\ell-1} \lambda^{-\ell-k} \end{aligned}$$

Clearly, the second term cancels the contribution of the 0 singularity, with the net result of

$$\frac{1}{2\pi i} \oint (y-\lambda)^{-\ell} \left(y^{k-1} + y^{-k-1} \right) dy = \binom{k-1}{\ell-1} \lambda^{k-\ell} \tag{5}$$

The RHS integral has only the 0 singularity contributing (the $y = \frac{1}{\lambda}$ singularity is outside the unit circle), yielding

$$\left(-\frac{1}{\lambda}\right)^{\ell} \cdot \frac{1}{(k-\ell)!} \cdot \left.\frac{d^k (y-\frac{1}{\lambda})^{-\ell}}{dy^{k-\ell}}\right|_{y=0} = \left(-\frac{1}{\lambda}\right)^{\ell} \cdot \frac{(-\ell)(-\ell-1)\dots(-k+1)}{(k-\ell)!} \cdot \left(-\frac{1}{\lambda}\right)^{-k} = \binom{k-1}{k-\ell} \lambda^{k-\ell}$$

which is the same as (5).

We should mention that in the most common case of $\ell = 1$, this answer boils down to λ^{k-1} (for $\ell = 2$ we are getting $(k-1)\lambda^{k-2}$, for $\ell = 3$ this becomes $\frac{(k-1)(k-2)}{2}\lambda^{k-3}$, etc.).

Example: Suppose

$$\omega(\beta) = \frac{12 + 6\cos(\beta) - 6\cos(2\beta) - 12\cos(3\beta)}{17 + 8\cos(2\beta)}$$

This is equivalent to

$$\omega(y) = -\frac{3}{2} - \frac{3}{4}\left(y + \frac{1}{y}\right) + \frac{\frac{15}{32} + \frac{5i}{8}}{y - \frac{i}{2}} + \frac{\frac{15}{32} - \frac{5i}{8}}{y + \frac{i}{2}} + \frac{\frac{15}{8} - \frac{5i}{2}}{y - 2i} + \frac{\frac{15}{8} + \frac{5i}{2}}{y + 2i}$$

where $y = \exp(i\beta)$. We thus get (almost immediately)

$$\rho_k = \left(\frac{15}{32} + \frac{5i}{8}\right) \left(\frac{i}{2}\right)^{k-1} + \left(\frac{15}{32} - \frac{5i}{8}\right) \left(-\frac{i}{2}\right)^{k-1} \qquad \text{for } k > 1$$

and

$$\rho_1 = \left(\frac{15}{32} + \frac{5i}{8}\right) + \left(\frac{15}{32} - \frac{5i}{8}\right) - \frac{3}{4} = \frac{3}{16}$$